

## The Dynamical Matrix

Recall: Hamiltonian:  $\mathcal{H} = \sum_l \frac{p_l^2}{2m_l} + \frac{1}{2} \sum_{l, l'} \psi(l, l') u_l u_{l'}$

Classical eqn's of motion:

$$1. \frac{\partial \mathcal{H}}{\partial p_l} = \dot{r}_l = \dot{u}_l \quad 2. \frac{\partial \mathcal{H}}{\partial R_l} = -\dot{p}_l = \frac{\partial \mathcal{H}}{\partial u_l}$$

$$1. \frac{\partial \mathcal{H}}{\partial p_l} = \dot{u}_l \Rightarrow p_l = m_l \dot{u}_l$$

$$2. \frac{\partial \mathcal{H}}{\partial R_l} = -\dot{p}_l = \frac{\partial \mathcal{H}}{\partial u_l} \Rightarrow \frac{\partial \mathcal{H}}{\partial u_l} = \frac{1}{2} \frac{\partial}{\partial u_l} \left\{ \sum_{l', l''} \psi(l', l'') u_{l'} u_{l''} \right\}$$

Note  $l, l', l''$   
are just indices

$$= \frac{1}{2} \left\{ \sum_{l'} \psi^0(l', l) u_{l'} + \sum_{l''} \psi^0(l, l'') u_{l''} \right\}$$

$$= \sum_{l'} \psi^0(l, l') u_{l'}$$

and:  $\dot{p}_l = m_l \ddot{u}_l$

$$\therefore m_l \ddot{u}_l = - \sum_{l'} \psi^0(l, l') u_{l'}$$

OR  $m_l \ddot{u}_l + \sum_{l'} \psi^0(l, l') u_{l'} = 0$

This is valid in general for an interaction potential energy function which is dependent on the ion positions.

$$M_L \ddot{u}_L + \sum_{L'} \psi^0(L, L') u_{L'} = 0$$

→ leads to set of  $N$  coupled diff. eqn's.

$$\text{e.g. } M_1 \ddot{u}_1 + \psi^0(1,1)u_1 + \psi^0(1,2)u_2 + \dots + \psi^0(1,N)u_N = 0$$

$$\vdots$$

$$M_N \ddot{u}_N + \psi^0(N,1)u_1 + \psi^0(N,2)u_2 + \dots + \psi^0(N,N)u_N = 0$$

Assume solutions of the form:

$$u_L(t) = u_L^0 e^{\pm i\omega t} \quad \rightarrow \text{amp. of oscillation}$$

$$\Rightarrow -M_L \omega^2 u_L^0 + \sum_{L'} \psi^0(L, L') u_{L'}^0 = 0$$

or

$$\sum_{L'} \left[ -M_L \omega^2 \delta_{LL'} + \psi^0(L, L') \right] u_{L'}^0 = 0$$

$N \times N$  matrix.  $\equiv L$

$$\text{define: } -L_{LL'} \equiv -M_L \omega^2 \delta_{LL'} + \psi^0(L, L')$$

$$\Rightarrow L u^0 = 0$$

$N \times N$  matrix  $\leftarrow$  vector of  $u_L^0$ 's

Recall:  $\psi^0(l, l') \equiv \frac{\partial^2 \omega}{\partial R_{l'} \partial R_l} \Big|_0 \rightarrow \text{real}$

$$\psi^0(l, l') = \psi^0(l', l)$$

$$L_{l, l'} = M_l \omega^2 \delta_{ll'} - \psi^0(l, l')$$

$\rightarrow$  elements are real & symmetric

$\therefore$  Matrix  $L$  is Hermitian

Our set of equations:  $\sum_{l'} L_{ll'} u_{l'}^0 = 0$

has (non trivial) solutions provided  $|L| = 0$

$\rightarrow$  Results in solution of  $N$  roots for  $\omega^2$

$\rightarrow$  Since  $L$  is Hermitian, the roots i.e.  $\omega^2$  are real

$\rightarrow$  while this corresponds to  $2N$  soln for  $\omega$   
( $\pm \sqrt{\omega^2}$ )

$$u_l(t) \propto e^{i\omega t} + e^{-i\omega t}$$

$\therefore$  still just  $N$  soln's.  $\rightarrow$  1D

For  $N$  degrees of freedom,  $N$  roots

i.e. 1D  $\rightarrow$   $N$  roots  
2D  $\rightarrow$   $2N$   
3D  $\rightarrow$   $3N$

→ each value of  $\omega^2$  represents a phonon frequency or normal mode of vibration

Assume solutions:

$$U_l^0(k) = A(k) \frac{e^{ik \cdot R_l^0}}{\sqrt{M_l}}$$

amplitude of mode  $k$  (pointing to  $A(k)$ )  
 normalization (pointing to  $\frac{1}{\sqrt{M_l}}$ )  
 in 1D  $R_l^0 = a l$ . (pointing to  $R_l^0$ )

→ apply periodic boundary conditions:  $U_L = U_{L+N}$

$$k = \frac{2\pi}{aN} m, \quad m \text{ integer}$$

→ for 1D but sim. expressions for 2 → 3D.

Recall, there are  $N$  modes, i.e. # of  $m$ 's. By convention, limit  $m$  to:

$$m \in \left( -\frac{N}{2} + 1, \frac{N}{2} \right) \Rightarrow \text{total } N \text{ modes}$$

Verify solutions

$$\sum_{l'} L_{ll'} U_{l'}^0 = 0 = \sum_{l'} \left( M_l \omega^2(k) \delta_{ll'} - \psi^0(l, l') \right) A(k) \frac{e^{ikR_{l'}^0}}{\sqrt{M_{l'}}$$

$$= M_l \omega^2(k) A(k) \frac{e^{ikR_l^0}}{\sqrt{M_l}} - \sum_{l'} \frac{e^{ikR_{l'}^0}}{\sqrt{M_{l'}}} A(k) \psi^0(l, l') e^{ik(R_{l'}^0 - R_l^0)}$$

= 1 when  $l' = l$

$$0 = \sqrt{M_l} A(k) e^{ikR_l^0} \left[ \omega^2(k) - \sum_{l'} \frac{\psi^0(l, l')}{\sqrt{M_l M_{l'}}} e^{ik(R_{l'}^0 - R_l^0)} \right]$$

$\therefore$

$$= 0$$

$$\therefore \omega^2(k) = \sum_{l'} \frac{\psi^0(l, l')}{\sqrt{M_l M_{l'}}} e^{ik(R_{l'}^0 - R_l^0)}$$

$\equiv D(k) \rightarrow$  Dynamical Matrix.

$$\omega^2(k) = D(k)$$

$$D(k) \equiv \sum_{l'} \frac{\psi^0(l, l')}{\sqrt{M_l M_{l'}}} e^{ik(R_{l'}^0 - R_l^0)}$$

Note that  $D(k)$  does not depend on  $l$ :

$\rightarrow$  take  $M_{l'} = M_l = M$

$\rightarrow R_l^0$  are lattice vectors  $\therefore R_{l'}^0 - R_l^0 \equiv R_{l'-l}^0$  is a lattice vector

$\rightarrow \psi^0(l, l')$  depends on separation of  $l$  &  $l'$

let  $l' - l \equiv m$

$$D(k) = \frac{1}{M} \sum_m e^{ikR_m^0} \psi^0(m)$$

$$D(k) = \frac{1}{M} \sum_m e^{ikR_m^0} \varphi^0(m) = \omega^2(k)$$

generalization of:

$$\omega^2 = \frac{k}{m}$$

mass

geometric  
phase factor

Force constants

One more note:

$$\sum_{l, l'} \varphi^0(l, l') = 0$$

→ think Newton's 3<sup>rd</sup> Law.

$$\text{i.e. } \sum_m \varphi^0(m) = 0$$

$$\Rightarrow \varphi^0(0) + \sum_{m \neq 0} \varphi^0(m) = 0$$

↳ separate self interaction

$$\therefore D(k) = \frac{1}{M} \sum_{m \neq 0} e^{ikR_m^0} \varphi^0(m) + \frac{1}{M} \varphi^0(0) = -\sum_{m \neq 0} \varphi^0(m)$$

$$D(k) = \frac{1}{M} \sum_{m \neq 0} (e^{ikR_m^0} - 1) \varphi^0(m)$$

geometric  
factor depends  
only on lattice

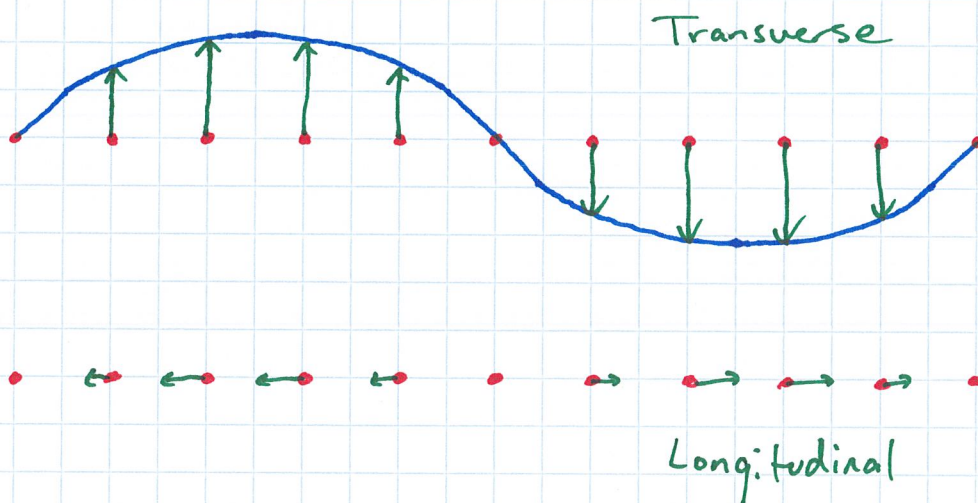
"self-force" components  
removed

## Physical Interpretation of Solutions

$$u_{\ell}(k, t) = \frac{E(k)}{\sqrt{M}} e^{i k R_{\ell}^0} e^{-i \omega(k) t}$$

↳ displacement of ion  $\ell$  vibrating in mode  $k$

- Normal modes are collective vibrations.



- sound waves are low frequency longitudinal & trans. phonons (20 Hz - 20 KHz)

$$\text{largest } \lambda ? : k = \frac{2\pi}{\lambda} = \frac{2\pi m}{Na} \Rightarrow m=1 \Rightarrow \lambda = Na$$

size of crystal

$$\text{smallest } \lambda ? : m = \frac{N}{2} \Rightarrow \lambda = 2a$$

↳ 2 unit cells